# LASER OPTOGALVANIC WAVELENGTH CALIBRATION WITH A COMMERCIAL HOLLOW CATHODE IRON-NEON DISCHARGE LAMP 

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#### Abstract

-351 optogalvanic transitions have been observed in the $337-598 \mathrm{~nm}$ wavelength region using an iron-neon hollow cathode discharge lamp and a pulsed tunable dye laser. 223 of these have been identified as transitions associated with neon energy levels. These optogalvanic transitions have allowed, in conjunction with interference fringes recorded concomitantly with an etalon, the calibration of the dye laser wavelength with $0.3 \mathrm{~cm}^{-1}$ accuracy.


## INTRODUCTION

Wavelength calibration of tunable lasers is important in experimental laser spectroscopy. An iodine absorption cell can be used for precise calibration of the spectral region $500-675 \mathrm{~nm}$ in the visible, since a good atlas exists for that region. ${ }^{1}$ However, in the blue and near u.v. regions of the electromagnetic spectrum, there is a paucity of suitable species for wavelength calibration of tunable lasers. Tellurium has been employed by Miller. ${ }^{2}$ However, its spectral atlas is not easily accessible. The optogalvanic (OG) effect provides a good solution for wavelength calibration inadequacies in the visible and near u.v. regions of the electromagnetic spectrum.
The OG effect, discovered by Penning ${ }^{3}$ in 1928, is the impedance change of a neon discharge tube when illuminated by a second neon discharge tube. In 1976, Green et al ${ }^{4}$ observed the OG effect with a tunable dye laser and measured the changes in voltage across a low-pressure gas discharge tube irradiated by a laser beam that was tuned to the wavelength of transition of a species present in the discharge. These authors reported optogalvanic signals for transitions involving lithium, sodium, calcium, barium, uranium, neon, and helium. It was pointed out that the OG effect might have significant impact in optical spectroscopy and other areas of applied physics and chemistry. Extensive work has since been carried out to understand the OG effect in some detail and to use it in laser spectroscopy. Precise wavelength calibration of tunable lasers is an important application of the OG effect.

OG lines have been observed in case of various species sputtered from hollow cathodes ${ }^{4-9}$ and from fill gases such as neon and argon. ${ }^{4.10-19}$ A larger discharge current ( $>20 \mathrm{~mA}$ ) is needed for exciting OG transitions associated with species sputtered from cathodes as compared to the signal from the gas atoms. Uranium, for example, is a good candidate for producing the OG effect, because it is an excellent wavelength standard. In addition, uranium gives rise to spectral lines that are narrow and do not possess hyperfine structure over a wide spectral region extending from the infrared to the ultraviolet. However, the OG signals arising from discharge gases are usually much stronger than those from cathode elements. OG transitions from argon have been observed in the wavelength regions $367-422 \mathrm{~nm},{ }^{19} 415-670 \mathrm{~nm},{ }^{15} 425-700 \mathrm{~nm},{ }^{14} 420-740 \mathrm{~nm},{ }^{16} 555-575 \mathrm{~nm},{ }^{17}$ $727-772 \mathrm{~nm},{ }^{13} 360-740 \mathrm{~nm},{ }^{18} 2440-2780 \mathrm{~nm}$. ${ }^{11}$ OG signals with commercial neon lamps, which are
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extensively used as spectral light sources, have been reported in the $540-750 \mathrm{~nm} .16 .13 .20$ and $2440-2780 \mathrm{~nm}^{11}$ regions. However, neon transitions do not seem to have been reported in the blue and near u.v. regions. It has been pointed out ${ }^{14}$ that it is difficult to obtain OG resonances with neon in wavelength regimes shorter than 580 nm owing to low oscillator strengths of the resonances.

In a series of measurements we have recently observed laser-assisted OG signals with neon in the $337-598 \mathrm{~nm}$ region by employing a commercial iron-neon hollow cathode lamp. 223 of these OG transitions have been identified to be associated with neon energy levels. Surprisingly, we have found that there are more OG lines of neon in the near-u.v. region than in the yellow and red regions, and interestingly the $O G$ signals in wavelength regimes shorter than 360 nm are very strong. Interference fringes from two surfaces of an etalon were also recorded simultaneously to determine an unknown transition wavelength between two neon transitions. The optogalvanic lines reported here should be useful for accurate wavelength calibration of tunable lasers in high-resolution laser spectroscopy.

## EXPERIMENTAL

Figure 1 shows the experimental arrangement for the $O G$ wavelength calibration in our laser-induced fluorescence investigations. A dye laser (DL) is pumped by an excimer laser (EL) running at 10 Hz . The output beam has a pulse duration of about 20 nsec . and a nominal linewidth of $0.07 \mathrm{~cm}{ }^{1}$ without any intracavity etalon. The tuning range $336-600 \mathrm{~nm}$ was covered by the following Exciton laser dyes: P-Terphenyl, TMQ, BPBD, Exalite 376, Exalite 389, Exalite 398. DPS, Bis-MSB, Coumarin 440, Coumarin 460, Coumarin 480, Coumarin 500, Coumarin 540 and Rhodamine 590 . An uncoated quartz wedge (WD) was inserted in the optical path of the primary beam to pick off two weak beams (each about $5 \%$ of the primary pulse energy). One of the beams (unfocussed and of typical pulse energy $100 \mu \mathrm{~J}$ ) enters the cathode of a commercial iron-neon hollow cathode lamp through a 1 mm diameter aperture. The second beam traverses a negative lens (NL) and illuminates an uncoated, parallel-faced 6 mm thick quartz disk at a small angle of incidence ( $1-2 \mathrm{deg}$ ) which serves as a low-finesse etalon (ET). The interference pattern, generated by the reflection beams from the front and rear surfaces of the disk, is recorded after passage through a pinhole aperture (AP) by a photodiode (PD).


Fig. 1. Experimental arrangement for laser optogalvanic wavelength calibration. $\mathrm{EL}=$ Excimer Laser, $D L=D y e \quad$ Laser $\quad W D=W e d g e . \quad P D=$ Photodiode $\quad \quad A P=$ Aperture,$\quad N L=$ Negative Lens, $\mathrm{OSC}=$ Oscilloscope, $\mathrm{HL}=$ Hollow Cathode Discharge Lamp, $\mathrm{PC}=$ Personal Computer, $\mathrm{BC}=\mathrm{Boxar}$ Averager, $C=$ Capacitor, $R=$ Resistor, $P S=$ Power Supply.

A high-voltage power supply (PS) and a ballast resistor (R) of $20 \mathrm{k} \Omega$ were used for the iron-neon lamp. The discharge current was set at 1 mA and the voltage across the lamp was 160 V . When the laser pulse is resonantly absorbed by the discharge medium, the voltage across the lamp varies, and these variations are coupled via a $0.05 \mu \mathrm{~F}$ capacitor to a boxcar integrator (BC). The temporal evolution of the signal was recorded by a digital oscilloscope. The outputs of the boxcar and the photodiode were recorded with a microcomputer-aided data acquisition system. The OG signal and the interference fringes were recorded simultaneously as illustrated in Fig. 2.

## RESULTS AND DISCUSSION

351 OG transitions have been recorded in the $336.99-597.55 \mathrm{~nm}$ wavelength region and are listed in Table 1.223 of these lines have been identified as due to neon. The assignments of the transitions and the associated emission intensities follow Ref. 21. The OG transitions have been identified according to the J-L coupling scheme of Racah. ${ }^{22}$ The electronic configuration of the neon ion is $1 s^{2} 2 s^{2} 2 p^{5}$ with the spectral terms ${ }^{2} P_{1.5}^{0}$ and ${ }^{2} P_{0.5}^{0}$. When a $3 s$ electron is added to the ${ }^{2} P_{1.5}^{0}$ configuration, the generated energy levels are designated $3 s[1.5]^{\circ} 1$ and $3 s[1.5]^{\circ} 2$. The number 1.5 within the parentheses represents one of the vector sums of the ionic $J$ and the $l$ value of the added electron ( $l=0$ for an $s$ electron) and the superscript 0 indicates odd parity. The number on the right of the parenthesis is the total $J$ value; it is the vectorial sum of the ionic $J$ and the spin of the added $3 s$ electron. Similarly, if a $3 p$ electron is added to the ionic term ${ }^{2} P_{1.5}^{0}$, the generated energy levels are $3 p^{\prime}[0.5] 0,3 p^{\prime}[0.5] 1,3 p^{\prime}[1.5] 1$ and $3 p^{\prime}[1.5] 2$, where the prime indicates that the level is generated from the $J=0.5$ ionic level. 128 OG lines have been recorded and tabulated in Table 1, but have not been assigned. These are confined mainly to the spectral region $361.5-427.6 \mathrm{~nm}$. These transitions are probably due to weak emission lines not listed in the neon atlas. Because the current in the discharge lamp was only 1 mA , we did not expect to observe any iron transitions. The time evolution of the OG signals was recorded by a digital oscilloscope. The amplitudes of the signals were varied from 10 mV to 2 V and weaker lines were not recorded. The waveforms had different shapes for the various kinds of observed transitions and depended on the magnitude of the discharge current. Most observed pulses had an initial negative peak, rose to cross the base line and became positive. Figure 3(a), (b), and (c) shows the OG signal waveforms for three distinct transitions at wavelengths (in air) $3510.721,3515.190$, and $3520.471 \AA$, respectively. Each curve was an average over 1000 events. All the three waveforms were recorded with the discharge current maintained at 0.5 mA and the laser pulse energy at $20 \mu \mathrm{~J}$. However, it can be


Fig. 2. Sample optogalvanic spectrum of neon (top) and the interference pattern of the etalon (bottom) around 350 nm .

Table 1. Optogalvanic neon transitions in the $3369-5980 \AA$ region.

| Wavelength in $\operatorname{Air}(\dot{\AA})$ | Intensity of Emission | Ne Transition | OG Signal Intensity(mV) | Laser Energy <br> ( $\mu \mathrm{J}$ ) |
| :---: | :---: | :---: | :---: | :---: |
| 3369.9069 | 700 | 3s $[1.5]^{0} 2-4 p^{\prime}[0.5] 1$ | 1500 | 50 |
| 3375.6489 | 50 | 3s [1.5] ${ }^{0} 2-4 p^{\prime}[1.5] 1$ | 100 | 55 |
| 3417.9031 | 500 | 3s [1.5] ${ }^{0} 1-4 p^{2}[1.5] 2$ | 700 | 70 |
| 3418.007 | 50 | $3 \mathrm{~s}[1.5]^{0} 1-4 p^{\prime}[0.5] 1$ | 400 | 60 |
| 3423.9120 | 50 | 3s [1.5] ${ }^{0} 1-4 p^{2}[1.5] 1$ | 200 | 60 |
| 3447.7022 | 200 | $3 \mathrm{~s}[1.5]^{0} 2-4 \mathrm{p}[1.5] 2$ | 20 | 45 |
| 3450.7641 | 50 | 3s [1.5] ${ }^{0}$ - $4 \mathrm{p}[1.5] 1$ | 1300 | 40 |
| 3454.1942 | 100 | 3s [1.5] ${ }^{0} 1-4 \mathrm{p}[0.5] 0$ | 600 | 20 |
| 3460.5235 | 100 | $3 s^{\prime}[0.5]^{\circ} 0-4 p^{\prime}(0.5] 1$ | 600 | 25 |
| 3464.3385 | 100 | 3s [1.5] ${ }^{\circ} 2-4 \mathrm{p}[2.5] 2$ | 1000 | 25 |
| 3466.5781 | 200 | $3 s^{\prime}[0.5]^{0} 0-4 p^{\prime}[1.5] 1$ | 1000 | 30 |
| 3472.5706 | 500 | $3 \mathrm{~s}[1.5]^{0} 2-4 \mathrm{p}[2.5] 3$ | 300 | 30 |
| 3498.0632 | 100 | 3s [1.5] ${ }^{0} 1-4 \mathrm{p}[1.5] 2$ | 150 | 35 |
| 3501.2154 | 200 | 3s [1.5] ${ }^{0} 1-4 \mathrm{p}[1.5] 1$ | 200 | 35 |
| 3510.7207 | 50 | 3s [1.5] ${ }^{\circ} 2-4 \mathrm{p}[0.5] 1$ | 800 | 40 |
| 3515.1900 | 200 | 3s [1.5] ${ }^{0} 1-4 \mathrm{p}[2.5] 2$ | 150 | 40 |
| 3520.4714 | 1000 | $3 s^{\prime}[0.5]^{0} 1-4 p^{\prime}[0.5] 0$ | 50 | 45 |
| 3593.5263 | 500 | $\left.3 s^{\prime}[0.5]^{0} 1-4 \mathrm{p}{ }^{\prime} 1.5\right] 2$ | 25 | 50 |
| 3600.1694 | 100 | $3 s^{\prime}[0.5]^{0} 1-4 p^{\prime}[1.5] 1$ | 15 | 45 |
| 3609.1787 | 50 | $3 s^{\prime}[0.5]^{0} 0-4 p[0.5] 1$ | 150 | 50 |
| 3614.23 |  |  | 10 | 50 |
| 3621.87 |  |  | 5 | 55 |
| 3626.19 |  |  | 10 | 55 |
| 3628.82 |  |  | 25 | 60 |
| 3633.6643 | 100 | $3 s^{\prime}[0.5]^{0} 1-4 p[0.5] 0$ | 35 | 80 |
| 3682.2421 | 100 | $3 s^{\prime}[0.5]^{0} 1-4 p[1.5] 2$ | 25 | 50 |
| 3685.02 |  |  | 22 | 50 |
| 3685.7351 | 100 | $3 s^{\prime}[0.5]^{0} 1-4 p[1.5] 1$ | 28 | 50 |
| 3701.2247 | 40 | $3 s^{\prime}[0.5]^{\circ} 1-4 \mathrm{p}[2.5] 2^{\prime}$ | 35 | 70 |
| 3754.2148 | 50 | $3 \mathrm{~s}^{\prime}[0.5]^{0} 1-4 \mathrm{p}[0.5] 1$ | 18 | 60 |
| 3829.77 |  |  | 10 | 45 |
| 3862.57 |  |  | 15 | 120 |
| 3876.00 |  |  | 15 | 130 |
| 3880.39 |  |  | 5 | 130 |
| 3893.02 |  |  | 20 | 130 |
| 3915.21 |  |  | 20 | 110 |
| 3922.25 |  |  | 6 | 90 |
| 3922.96 |  |  | 7 | 90 |
| 3924.57 |  |  | 7 | 90 |
| 3927.18 |  |  | 8 | 80 |
| 3927.31 |  |  | 8 | 80 |
| 3933.58 |  |  | 8 | 80 |
| 3942.19 |  |  | 18 | 80 |
| 3943.33 |  |  | 34 | 60 |
| 3946.26 |  |  | 20 | 100 |
| 3947.98 |  |  | 14 | 100 |
| 3952.70 |  |  | 22 | 100 |
| 3953.58 |  |  | 14 | 100 |
| 3954.00 |  |  | 10 | 100 |
| 3954.72 |  |  | 16 | 100 |
| 3960.45 |  |  | 22 | 110 |
| 3962.92 |  |  | 18 | 110 |
| 3969.66 |  |  | 24 | 120 |
| 3969.78 |  |  | 25 | 120 |
| 3972.59 |  |  | 9 | 120 |
| 3972.80 |  |  | 20 | 120 |
| 3981.16 |  |  | 23 | 130 |
| 3981.33 |  |  | 23 | 130 |
| 3984.065 | 2 | 3p [0.5] $1-9 \mathrm{~d}$ ' 11.5$]^{0} 1$ | 22 | 130 |
| 3984.253 | 7 | $3 \mathrm{p}[0.5] 1-9 \mathrm{~d}^{\prime}[2.5]^{0} 2$ | 35 | 130 |
| 3985.12 |  |  | 12 | 130 |
| 3995.50 |  |  | 28 | 120 |
| 3995.721 | 1 | 3p $[0.5] 1-13 \mathrm{~d}[0.5]^{0} 1$ | 27 | 120 |
| 3998.594 | 1 | 3p $[0.5] 1-10 s^{\prime}[0.5]^{0} 1$ | 34 | 120 |
| 3999.263 | 1 | 3p $[0.5] 1-10 s^{\prime}[0.5]^{0} 0$ | 14 | 120 |
| 4000.24 |  |  | 14 | 120 |
| 4000.44 |  |  | 24 | 120 |
| 4013.752 | 1 | 3p [0.5] $1-12 \mathrm{~d}[1.5]^{0} 2$ | 28 | 120 |
| 4013.995 | 2 | $3 \mathrm{p}[0.5] 1-12 \mathrm{~s}[0.5]^{0} 1$ | 25 | 110 |
| 4020.015 | 2 | $3 \mathrm{p}[0.5] 1-13 \mathrm{~s}[1.5]^{0} 2$ | 20 | 100 |
| 4037.262 | 5 | 3p [0.5] 1-11d [1.5] ${ }^{0} 2$ | 20 | 100 |
|  |  |  | continued oppesite |  |

Table l-continued.

| Wavelength in $\operatorname{Air}(\mathrm{A})$ | Intensity of Emission | Ne Transition | OG Signal Intensity (mV) | Laser Energy ( $\mu \mathrm{J}$ ) |
| :---: | :---: | :---: | :---: | :---: |
| 4037.615 | 15 | 3p [0.5] 1-11d [0.5] ${ }^{0} 1$ | 20 | 100 |
| 4037.696 | 5 | 3p $[0.5] 1-11 \mathrm{~d}[0.5]^{0} 0$ | 10 | 100 |
| 4042.327 | 10 | $3 \mathrm{p}[0.5] 1-8 \mathrm{~d}^{\prime}(1.5)^{\circ} 1$ | 20 | 110 |
| 4042.642 | 50 | 3p [0.5] 1-8d'[2.5] ${ }^{0} 2$ | 25 | 110 |
| 4045.04 |  |  | 10 | 110 |
| 4045.662 | 2 | 3p [0.5] 1-12s [1.5] ${ }^{\circ} 2$ | 12 | 110 |
| 4064.036 | 50 | 3p [0.5] 1-9s' 0.5$]^{0} 1$ | 15 | 110 |
| 4064.829 | 15 | 3p $[0.5] 1-9 s^{\prime}[0.5]^{0} 0$ | 10 | 120 |
| 4068.835 | 30 | 3p [0.5] 1-10d [1.5p ${ }^{2}$ | 20 | 120 |
| 4069.243 | 30 | 3p [0.5] 1-10d [0.5] ${ }^{0} 1$ | 20 | 120 |
| 4069.389 | 5 | 3p [0.5] 1-10d [0.5] ${ }^{0} 0$ | 13 | 120 |
| 4079.359 | 2 | 3p $[0.5] 1-11 \mathrm{~s}[1.5]^{0} 1$ | 9 | 110 |
| 4080.148 | 50 | 3p $[0.5] 1-11 \mathrm{~s}[1.5]^{\circ} 2$ | 15 | 110 |
| 4087.82 |  |  | 8 | 110 |
| 4097.79 |  |  | 8 | 100 |
| 4111.84 |  |  | 10 | 90 |
| 4112.100 | 15 | 3p [0.5] 1 - 9d [1.5] ${ }^{0} 2$ | 12 | 80 |
| 4112.694 | 20 | 3p [0.5] 1-9d [0.5] ${ }^{0} 1$ | 12 | 80 |
| 4112.885 | 10 | 3p [0.5] 1 - 9d [0.5] ${ }^{\circ}$ | 8 | 80 |
| 4125.14 |  |  | 8 | 70 |
| 4126.941 | 2 | 3p [0.5] $1-10 \mathrm{~s}[1.5]^{\circ} 1$ | 12 | 50 |
| 4128.072 | 30 | 3p [0.5] $1-10 \mathrm{~s}[1.5]^{\circ} 2$ | 24 | 50 |
| 4130.512 | 20 | 3p [0.5] 1-7d'[1.5] ${ }^{0} 1$ | 15 | 60 |
| 4131.054 | 70 | 3p [0.5] 1-7d'[2.5] ${ }^{\circ} 2$ | 20 | 70 |
| 4144.51 |  |  | 17 | 90 |
| 4148.81 |  |  | 10 | 100 |
| 4149.50 |  |  | 10 | 100 |
| 4151.37 |  |  | 15 | 100 |
| 4152.18 |  |  | 12 | 100 |
| 4154.31 |  |  | 25 | 110 |
| 4155.22 |  |  | 20 | 110 |
| 4157.65 |  |  | 30 | 110 |
| 4158.70 |  |  | 22 | 110 |
| 4160.82 |  |  | 10 | 110 |
| 4161.51 |  |  | 30 | 110 |
| 4162.69 |  |  | 26 | 110 |
| 4164.802 | 50 | 3p $[0.5] 1-8 s^{\prime}[0.5]^{0} 1$ | 20 | 110 |
| 4166.091 | 30 | $3 \mathrm{p}[0.5] 1-8 s^{\prime}[0.5]^{0} 0$ | 40 | 120 |
| 4167.32 |  |  | 30 | 120 |
| 4169.69 |  |  | 20 | 120 |
| 4171.08 |  |  | 40 | 120 |
| 4171.74 |  |  | 10 | 120 |
| 4172.74 |  |  | 35 | 120 |
| 4173.966 | 2 | 3p [0.5] 1-8d [1.5] ${ }^{0} 1$ | 15 | 120 |
| 4174.369 | 70 | 3p [0.5] 1-8d [1.5] ${ }^{\circ}$ | 28 | 120 |
| 4175.223 | 60 | 3p $[0.5] 1-8 \mathrm{~d}[0.5]^{\circ} 1$ | 22 | 120 |
| 4175.488 | 40 | $3 \mathrm{p}[0.5] 1-8 \mathrm{~d}[0.5]^{\circ} 0$ | 20 | 120 |
| 4176.44 |  |  | 10 | 130 |
| 4177.16 |  |  | 40 | 130 |
| 4179.11 |  |  | 35 | 130 |
| 4180.37 |  |  | 10 | 130 |
| 4182.03 |  |  | 10 | 130 |
| 4183.39 |  |  | 10 | 130 |
| 4184.36 |  |  | 45 | 130 |
| 4186.75 |  |  | 45 | 130 |
| 4190.65 |  |  | 15 | 130 |
| 4191.80 |  |  | 10 | 130 |
| 4193.00 |  |  | 50 | 130 |
| 4195.15 |  |  | 15 | 140 |
| 4195.78 |  |  | 40 | 140 |
| 4196.415 | 15 | 3p [0.5] $1-9 \mathrm{~s}[1.5]^{0} 1$ | 25 | 140 |
| 4198.099 | 70 | 3p $[0.5] 1-9 \mathrm{~s}[1.5]^{\circ} 2$ | 22 | 140 |
| 4200.40 |  |  | 18 | 140 |
| 4201.95 |  |  | 10 | 140 |
| 4203.270 | 2 | 3p [2.5]2-10d'(2.5] ${ }^{\circ} 3$ | 60 | 140 |
| 4203.43 |  |  | 50 | 140 |
| 4206.43 |  |  | 16 | 140 |
| 4206.83 |  |  | 40 | 140 |

Table l-continued.

| Wavelength in $\operatorname{Air}(\AA)$ | Intensity of Emission | Ne Transition | OG Signal Intensity(mV) | Laser Energy ( $\mu \mathrm{J}$ ) |
| :---: | :---: | :---: | :---: | :---: |
| 4208.30 |  |  | 10 | 140 |
| 4208.47 |  |  | 10 | 140 |
| 4210.58 |  |  | 10 | 140 |
| 4213.80 |  |  | 20 | 140 |
| 4214.87 |  |  | 10 | 140 |
| 4216.04 |  |  | 50 | 140 |
| 4216.24 |  |  | 55 | 140 |
| 4220.61 |  |  | 40 | 140 |
| 4222.58 |  |  | 20 | 140 |
| 4225.31 |  |  | 14 | 140 |
| 4230.78 |  |  | 15 | 140 |
| 4232.07 |  |  | 40 | 140 |
| 4232.33 |  |  | 50 | 140 |
| 4233.17 |  |  | 32 | 140 |
| 4236.47 |  |  | 15 | 140 |
| 4236.70 |  |  | 10 | 140 |
| 4237.80 |  |  | 40 | 130 |
| 4246.24 |  |  | 20 | 130 |
| 4249.538 | 2 | $3 p[2.5] 2-9 d^{\prime}[2.5]^{\circ} 2$ | 20 | 130 |
| 4250.28 |  |  | 15 | 130 |
| 4252.418 | 2 | $3 \mathrm{p}[2.5] 3-12 \mathrm{~d}[2.5]^{0} 3$ | 45 | 130 |
| 4252.775 | 2 | 3p [2.5] 3-12d [3.5] ${ }^{0} 4$ | 50 | 130 |
| 4256.498 | 2 | $3 \mathrm{p}[1.5] 1-10 \mathrm{~d}^{\prime}[2.5]^{\circ} 2$ | 12 | 130 |
| 4256.99 |  |  | 13 | 120 |
| 4259.40 |  |  | 10 | 120 |
| 4259.77 |  |  | 40 | 120 |
| 4261.78 |  |  | 10 | 120 |
| 4262.23 |  |  | 15 | 120 |
| 4262.479 | 2 | $3 \mathrm{p}[2.5] 2-13 \mathrm{~d}[3.5]^{\circ} 3$ | 20 | 120 |
| 4265.89 |  |  | 10 | 120 |
| 4267.724 | 5 | $3 \mathrm{p}[0.5] 1-7 \mathrm{~d}[1.5]^{0} 1$ | 20 | 120 |
| 4268.009 | 70 | 3p [0.5] $1.7 \mathrm{~d}[1.5]^{0} 2$ | 35 | 120 |
| 4268.30 |  |  | 15 | 120 |
| 4269.724 | 70 | 3p [0.5] 1 - $7 \mathrm{~d}[0.5]^{0} \mathrm{t}$ | 20 | 120 |
| 4270.267 | 50 | $3 \mathrm{p}[0.5] 1-7 \mathrm{~d}[0.5]^{\circ} 0$ | 15 | 120 |
| 4271.19 |  |  | 12 | 110 |
| 4274.656 | 50 | 3p [0.5] $1-6 d^{\prime}[1.5]^{0} 1$ | 25 | 110 |
| 4275.5598 | 70 | $3 \mathrm{p}[0.5] 1-6 d^{\prime}[1.5]^{0} 2$ | 25 | 110 |
| 4278.850 | 5 | $3 \mathrm{p}[2.5] 3-11 \mathrm{~d}[2.5]^{0} 3$,2 | 45 | 110 |
| 4279.06 |  |  | 30 | 110 |
| 4279.279 | 15 | $3 \mathrm{p}[2.5] 3-11 \mathrm{~d}[3.5]^{0} 4,3$ | 45 | 110 |
| 4282.84 |  |  | 12 | 110 |
| 4283.242 | 10 | 3p 12.5$] 2-12 \mathrm{~d}[3.5]^{\circ} 3$ | 20 | 100 |
| 4287.03 |  |  | 10 | 100 |
| 4288.541 | 5 | $3 \mathrm{p}[2.5]^{3}-12 \mathrm{~s}(1.5]^{\circ} 2$ | 30 | 100 |
| 4288.93 |  |  | 10 | 90 |
| 4289.94 |  |  | 12 | 90 |
| 4290.40 |  |  | 10 | 90 |
| 4291.976 | 2 | $\left.3 \mathrm{p} \mid 1.5] 2-10 d^{\prime} \mid 2.5\right]^{\circ} 2,3$ | 20 | 90 |
| 4292.46 |  |  | 12 | 80 |
| 4293.09 |  |  | 10 | 80 |
| 4303.02 |  |  | 10 | 70 |
| 4303.955 | 5 | $3 \mathrm{p}[1.5] 1-9 \mathrm{~d}^{\prime}[2.5]^{\circ} 2$ | 10 | 70 |
| 4306.2625 | 70 | 3p $[0.5] 1-8 \mathrm{~s}[1.5]^{0} 2$ | 20 | 100 |
| 4309.10 |  |  | 18 | 100 |
| 4310.130 | 2 | 3p $12.512-11 \mathrm{~d}[3.5]^{0} 3$ | 20 | 110 |
| 4314.110 | 1 | 3p [2.513-10d [2.5 $]^{\circ} 3$ | 45 | 110 |
| 4314.695 | 30 | $3 \mathrm{p}\left[2.513-10 \mathrm{~d}[3.5]^{\circ} 4,3\right.$ | 45 | 110 |
| 4316.008 | 15 | 3p [2.5]2-8d'[2.5] ${ }^{\circ} 3,2$ | 20 | 110 |
| 4327.265 | 10 | $3 \mathrm{p}[2.5] 3-11 \mathrm{~s}[1.5]^{0} 2$ | 34 | 120 |
| 4340.418 | 2 | $3 \mathrm{p}[2.5] 2-9 s^{\prime}[0.5]^{0} 1$ | 20 | 120 |
| 4346.036 | 15 | 3p [2.5] $2-10 \mathrm{~d}[3.5]^{\circ} 3$ | 19 | 130 |
| 4362.690 | 30 | $3 \mathrm{p}[2.5] 3-9 \mathrm{~d}[2.5]^{\circ} 3$ | 25 | 150 |
| 4363.228 | 2 | $3 \mathrm{p}[2.5] 3-9 \mathrm{~d}[1.5]^{\circ} 2$ | 25 | 150 |
| 4363.524 | 70 | $3 \mathrm{p}\left\{2.5 \mid 3-9 \mathrm{~d}[3.5]^{\circ} 4\right.$ | 40 | 150 |
| 4374.997 | 2 | $\left\{^{3 \mathrm{p}}[1.5] 2-12 \mathrm{~d}[2.5]^{\circ} 3\right.$ | 20 | 160 |
|  |  | $3 \mathrm{p}[1.5] 1-12 \mathrm{~s}[1.5]^{0} 1$ |  |  |
| 4381.220 | 30 | $3 \mathrm{p}[2.5] 3-10 \mathrm{~s}[1.5]^{\circ} 2$ | 35 | 160 |
| 4395.69 |  |  | 18 | 160 |
| 4422.5205 | 300 | $3 \mathrm{p}\|0.5\| 1-6 d[1.5]^{\circ} 2$ | 20 | 160 |

Table 1-continued.


Table I-continued.

| Wavelength in $\operatorname{Air}(\dot{\mathrm{A}})$ | Intensity of Emission | Ne Transition | OG Signal Intensity(mV) | Laser Energy <br> ( $\mu \mathrm{J}$ ) |
| :---: | :---: | :---: | :---: | :---: |
| 4837.3139 | 500 | $3 \mathrm{p}[0.5] 1-6 \mathrm{~s}[1.5]^{0} 2$ | 13 | 100 |
| 4842.941 | 50 | 3 p '10.5] 1-7d [1.5] ${ }^{0}$ | 8 | 100 |
| 4849.530 | 30 | $3 \mathrm{p} \cdot\left[\left.1.5\|2-8 \mathrm{~s}\| 1.5\right\|^{\circ} 2\right.$ | 8 | 100 |
| 4852.6571 | 100 | $3 p^{\prime}[0.5] 1-6 d^{\prime}[2.5]^{0} 2$ | 8 | 100 |
| 4859.604 | 15 | $3 \mathrm{p}[0.5] 0-8 \mathrm{~s}[1.5]^{0} 1$ | 8 | 90 |
| 4863.0800 | 100 | $3 \mathrm{p}[1.5] 2-6 \mathrm{~d}[2.5]^{0} 3$ | 13 | 90 |
| 4865.501 | 100 | $3 p[1.5] 2-6 \mathrm{~d}[1.5]^{0} 2$ | 11 | 90 |
| 4866.476 | 80 | $3 \mathrm{p}[1.5] 2$ - $6 \mathrm{~d}[3.5]^{0} 3$ | 10 | 90 |
| 4868.268 | 70 | $3 \mathrm{p}[1.5] 2$ - $6 \mathrm{~d}[0.5]^{0} 1$ | 8 | 90 |
| 4884.9170 | 1000 | $\left\{^{3} p^{\prime}[1.5] 2-7 s^{\prime} 0.5\right]^{0} 1$ | 19 | 80 |
|  |  | $3^{p} p[2.5] 2-5 d^{\prime}[2.5]^{9} 3$ |  | 80 |
| 4885.084 | 100 | $3 \mathrm{p}[2.5] 2-5 \mathrm{~d}^{\prime}[2.5]^{0} 2$ | 10 | 80 |
| 4905.19 |  |  | 8 | 70 |
| 4927.48 |  |  | 10 | 50 |
| 4928.235 | 70 | $3 \mathrm{p} \cdot[0.5] 1-7 \mathrm{~s}^{\prime}[0.5]^{0} 1$ | 7 | 50 |
| 4939.0457 | 100 | $3_{p},[1.5] 2-7 \mathrm{~s}[1.5]^{0} 1$ | 10 | 60 |
| 4944.9899 | 100 | $3 \mathrm{p}[1.5] 2-7 \mathrm{~s}[1.5]^{\circ} 2$ | 10 | 60 |
| 4955.382 | 150 | $3 \mathrm{p}[1.5] 1-5 \mathrm{~d}^{\prime}(1.5]^{\circ} 1$ | 11 | 60 |
| 4957.0335 | 1000 | $3 \mathrm{p}[1.5] 1-5 \mathrm{~d}^{\prime}[1.5]^{\circ} 2$ | 12 | 70 |
| 4973.538 | 100 | $3 \mathrm{p} \times[1.5] 1-6 \mathrm{~d}[2.5]^{\circ} 2$ | 8 | 70 |
| 4974.760 | 50 | $3 p^{\prime}[1.5] 1-6 d[1.5]^{0} 1$ | 6 | 70 |
| 4994.10 |  |  | 8 | 90 |
| 4994.930 | 150 | $3 p^{\prime}[1.5] 2-6 d[2.5]^{0} 3$ | 9 | 90 |
| 4997.482 | 15 | $3 p^{\prime}[1.5] 2-6 d[1.5]^{0} 2$ | 8 | 90 |
| 4997.96 |  |  | 12 | 90 |
| 4998.502 | 10 | $3 p^{\prime}\left[1.512-6 d[3.5]^{\circ} 3\right.$ | 12 | 90 |
| 5005.1587 | 500 | 3p [1.5]2-5d'[2.5] ${ }^{0} 3$ | 10 | 100 |
| 5022.870 | 25 | $\left.3 \mathrm{p}[2.5] 2-6 s^{\prime}[0.5]^{\circ}\right]$ | 8 | 100 |
| 5031.3504 | 250 | 3p [2.5]3-5d [2.5] ${ }^{0} 3$ | 23 | 100 |
| 5035.989 | 35 | 3p [2.5] 3 - $5 \mathrm{~d}[1.5]^{\circ} 2$ | 20 | 110 |
| 5037.7512 | 500 | 3p $[2.5] 3-5 \mathrm{~d}[3.5]^{\circ} 4$ | 25 | 110 |
| 5074.201 | 35 | $3 \mathrm{p}\left[2.512\right.$ - $5 \mathrm{~d}[2.5]^{\circ} 2$ | 12 | 110 |
| 5076.5816 | 35 | $3 \mathrm{p}[2.5] 2-5 \mathrm{~d} \mid 1.5]^{\circ} \mathrm{l}$ | 10 | 110 |
| 5078.762 | 15 | 3p $[2.5] 2-5 d[1.5]^{\circ} 2$ | 9 | 110 |
| 5080.3852 | 150 | $3 \mathrm{p}[2.5] 2-5 \mathrm{~d}[3.5]^{0} 3$ | 13 | 110 |
| 5083.968 | 25 | 3p $\left.[2.5] 2-5 d[0.5]^{\circ}\right]$ | 8 | 120 |
| 5113.6724 | 75 | $3 \mathrm{p}[0.5] 1-4 \mathrm{~d}[1.5]^{0} 1$ | 12 | 120 |
| 5116.5032 | 150 | 3p [0.5]1 - 4d' 1.5$]^{\mathrm{p}} 2$ | 14 | 120 |
| 5117.011 | 35 | $3 \mathrm{p}[0.5] 1-4 d^{\prime}[2.5]^{0} 2$ | 12 | 120 |
| 5122.257 | 150 | $3 \mathrm{p}^{\prime}[1.5] 1-5 \mathrm{~d}^{\prime}[1.5]^{\circ} 2$ | 9 | 120 |
| 5122.337 | 150 | $3 p^{\prime}\|1.5\| 1-5 d^{\prime}[2.5]^{0} 2$ | 9 | 120 |
| 5129 |  |  | 8 | 110 |
| 5144.9384 | 500 | $3 p^{\prime}[1.5] 2-5 d^{\prime}[2.5]^{6} 3$ | 13 | 110 |
| 5150.077 | 35 | $3 \mathrm{p}[1.5] 2-6 s^{\prime}[0.5]^{\circ} 1$ | 8 | 110 |
| 5151.9610 | 75 | 3p[1.5] 1-5d [2.5] ${ }^{\circ}$ | 8 | 110 |
| 5154.4271 | 50 | 3p [1.5] $1-5 \mathrm{~d} \mid 1.5]^{0} 1$ | 8 | 110 |
| 5156.667 | 50 | $3 \mathrm{p}[1.5] 1-5 \mathrm{~d}[1.5]^{\circ} 2$ | 8 | 110 |
| 5188.6122 | 150 | 3p[2.5]3-6s [1.5] ${ }^{\circ} 2$ | 15 | 100 |
| 5193.1302 | 150 | $3 p^{\prime}[0.5] 1-5 d^{\prime}\left[1.5 p^{\circ} 2\right.$ | 8 | 100 |
| 5193.2227 | 150 | $3 \mathrm{p} \cdot[0.5] 1-5 \mathrm{~d}^{\prime}(2.5]^{0} 2$ | 8 | 100 |
| 5203.8962 | 150 | 3p [1.5]2-5d [2.5] ${ }^{0} 3$ | 10 | 90 |
| 5208.8648 | 70 | 3p $[1.5] 2-5 \mathrm{~d}[1.5]^{\circ} 2$ | 10 | 90 |
| 5210.5672 | 50 | $3 \mathrm{p}[1.5] 2-5 \mathrm{~d}\left[3.50^{\circ} 3\right.$ | 10 | 90 |
| 5222.3517 | 50 | 3 p [0.5] $2-6 \mathrm{~s}[1.5]^{0} 1$ | 8 | 80 |
| 5326.3968 | 75 | 3p [0.5] 1 - $4 \mathrm{~d}[1.5]^{a} 1$ | 8 | 80 |
| 5330.7775 | 600 | 3p $[0.5] 1-4 d[1.5]^{0} 2$ | 10 | 80 |
| 5341.0938 | 1000 | 3p [0.5] $1-4 \mathrm{~d}[0.5]^{0} 1$ | 10 | 90 |
| 5343.2834 | 600 | $3 \mathrm{p}[0.5] 1$ - $4 \mathrm{~d}[0.5]^{\circ} 0$ | 8 | 90 |
| 5355.176 | 150 | $3 \mathrm{p}{ }^{\top}\|1.5\| 2-5 \mathrm{~d}[2.5]^{0} 3$ | 7 | 90 |
| 5398.49 |  |  | 7 | 100 |
| 5399.48 |  |  | 40 | 100 |
| 5400.5616 | 2000 | 3s $\left.[1.5]^{\circ} 1-3 \mathrm{p} \cdot 10.5\right] 0$ | 500 | 100 |
| 5433.6513 | 250 | $3 \mathrm{p}[0.5] 1-5 s^{\prime}(0.5]^{0} 1$ | 8 | 100 |
| 5448.5091 | 150 | 3p [0.5] $1-5 s^{\prime}[0.5]^{0} 0$ | 7 | 100 |
| 5511.485 | 15 | $3 \mathrm{p}[2.5] 3-4 \mathrm{~d}^{\prime}[2.5]^{0} 3$ | 8 | 100 |
| 5562.7662 | 500 | $3 \mathrm{p}[2.5] 2-4 \mathrm{~d}^{\prime}[2.5]^{0} 3$ | 8 | 90 |
| 5748.650 | 70 | 3p [2.5] $3-4 \mathrm{~d}[2.5]^{\circ} 2$ | 8 | 60 |
| 5764.063 | 3 | $3 \mathrm{p}[2.5] 3-4 \mathrm{~d}[3.5]^{0} 3$ | 8 | 70 |
| 5764.4188 | 700 | $3 \mathrm{p}[2.5] 3-4 \mathrm{~d}[3.5]^{0} 4$ | 10 | 70 |

Table 1-continued.

| Wavelength in $\operatorname{Air}(\AA)$ | Intensity of Emission | Ne Transition | 00 Signal Intensity(mV) | Laser Energy ( $\mu$ ) |
| :---: | :---: | :---: | :---: | :---: |
| 5770.307 | 50 | $3 p^{\prime}[0.5] 0-5 d^{\prime}[1.5]^{0} 1$ | 3 | 80 |
| 5804.4496 | 500 | 3p [2.5] $2-4 \mathrm{~d}[2.5]^{\circ} 2$ | 3 | 90 |
| 5820.1558 | 500 | 3p [2.5]2-4d [3.5] ${ }^{\circ} 3$ | 4 | 100 |
| 5881.8950 | 1000 | 3s [1.5] ${ }^{2}$ - 3p'10.5] 1 | 3800 | 90 |
| 5902.4623 | 50 | 3p'[1.5]2-4d'[2.5p3 | 4 | 80 |
| 5944.8342 | 500 | 3s [1.5] ${ }^{2}$ - 3p'[1.5] 2 | 2300 | 50 |
| 5965.4710 | 500 | 3p'[0.5] 1 - 4d'[1.5] ${ }^{\circ}$ | 4 | 40 |
| 5975.5340 | 600 | 3s [1.5] ${ }^{\circ} 2-3 \mathrm{p}$ [1.5] 1 | 1300 | 40 |

seen that the pulse shapes are quite different. The waveform in Fig. 3(a) corresponds to the neon transition $3 s[1.5]^{0} 2-4 p[0.5] 1$ at $3510.721 \AA$. It is a negative peak with a $17 \mu \mathrm{sec}$. full-width at half-maximum (FWHM) that rises and changes polarity to become a positive peak with a $100 \mu \mathrm{sec}$. FWHM. The waveform shown in Fig. 3(b) refers to the neon transition $3 s[1.5]^{0} 1-4 p[2.5] 2$ located at $3515.190 \AA$ that is oscillatory in nature. It has two negative peaks with FWHM of about 8 and $117 \mu \mathrm{sec}$. and a positive peak with a $13 \mu \mathrm{sec}$. FWHM in between. The waveform illustrated in Fig. 3(c) represents the neon transition $3 s^{\prime}[0.5]^{0} 1-4 p^{\prime}[0.5] 0$ at $3520.471 \AA$ and is an initially negative-going peak about $23 \mu \mathrm{sec}$. FWHM that decays gradually over $225 \mu \mathrm{sec}$. into more or less a flat tail. Figures 4(a) and (b) show two OG waveforms of the same neon transition $3 s[1.5]^{0} 2-4 p[0.5] 1$ at $3510.721 \AA$ recorded with two different discharge currents 1.00 and 0.25 mA respectively. The pulse shape shown in Fig. 4(b) was recorded with a discharge current of 0.25 mA and has a wider pulse-width of $19 \mu \mathrm{sec}$. FWHM compared to the pulse shown in Fig. 4(a) for 1.00 mA that had a pulse-width of $15 \mu \mathrm{sec}$. Despite the two waveforms having similar overall shapes their positive portions are different, in


Fig. 3. Sample waveforms of $O G$ signals with 0.5 mA discharge current and $20 \mu \mathrm{~J}$ laser pulse corresponding to neon transitions: (a) $3 s[1.5]^{0} 2-4 p[0.5] 1$ at $3510.721 \AA$, (b) $3 s[1.5]^{0} 1-4 p[2.5] 2$ at $3515.190 \AA$, and (c) $3 s^{\prime}[0.5]^{\circ} 1-4 p{ }^{\prime}[0.5] 0$ at $3520.471 \AA$.


Fig. 4. Waveforms of the OG signal from neon transition $3 x[1.5]^{\prime 2} 4 p[0.5]$ at $3510.721 \AA$ using a $20 \mu \mathrm{n}$. laser pulse and discharge current (a) 1.00 mA and (b) 0.25 mA .
that the higher current pulse is narrower, almost a third as wide as compared to the pulse recorded with the smaller discharge current. While recording the OG spectra special attention has to be paid in positioning the gate of the boxcar. The position of the gate changes the intensity and polarity of the signals. The intensities of the OG transitions listed in Table 1 correspond to the absolute amplitudes at the peak of the pulse.

A detailed discussion of the time evolution of the OG signals has been given by Smyth and Schenck, ${ }^{20}$ and by Reddy et al. ${ }^{23}$ Theoretical approaches to explain the OG phenomena have been carried out by Erez et al. ${ }^{24}$ van Veldhuizen et al, ${ }^{25}$ and Stewart et al. ${ }^{26}$ Our present investigations have been confined to a systematic and comprehensive identification and tabulation of laserinduced OG transitions in the u.v. and visible regions of the electromagnetic spectrum so as to calibrate the wavelengths of tunable lasers to within $0.3 \mathrm{~cm}{ }^{\prime}$ accuracy. No attempts have been made to explain the intensities and the temporal behavior of the OG signals.

The wavelength position between OG transitions can be determined by counting the simultaneously recorded interference fringes. The free spectral range of the etaton can be simply determined by dividing the wavenumber difference between OG transitions by the number of fringes in-between. For example, in Fig. 2, the wavenumber for position $P$ is given by

$$
\begin{equation*}
N_{\mathrm{P}}=N_{\mathrm{A}}+K_{\mathrm{PA}}\left(N_{\mathrm{B}}-N_{\mathrm{A}}\right) K_{\mathrm{BA}} \tag{1}
\end{equation*}
$$

where $N_{\mathrm{P}}, N_{\mathrm{A}}$ and $N_{\mathrm{B}}$ are the wavenumbers at positions $\mathrm{P}, \mathrm{A}$ and B , respectively, and $K_{\mathrm{P}}$ and $K_{\mathrm{B}}$, are the number of fringes between $A$ and $P$, and between $A$ and $B$, respectively. The nonlinearity of the refractive index of quartz, $\sim 10^{6} / \mathrm{nm}$ in the blue region, can be ignored because the linewidth of the OG signals was typically $0.3 \mathrm{~cm}{ }^{1}$ for the spectra recorded using an average over 10 shots for each data point. For instance, with $N_{\mathrm{A}}=28553.33 \mathrm{~cm}{ }^{\prime} . N_{\mathrm{B}}=28579.06 \mathrm{~cm}$ and $K_{\mathrm{BA}}=53.7$ and $K_{\mathrm{PA}_{\mathrm{A}}}=15.9$, the formula represented by Eq. (1) yields for the position P a value of $28560.95 \mathrm{~cm}^{\text {' }}$.

In summary, we have recorded 351 optogalvanic transitions using a commercial iron neon hollow cathode discharge lamp, and of which 223 lines have been assigned to neon transitions in the $337-598 \mathrm{~nm}$ region. To the best of our knowledge, most of these laser-assisted OG lines are being reported here for the first time and should be very useful for precise wavelength calibration.

An observation of the OG signal, along with simultaneous recording of interference fringes from an etalon, has provided precise and reliable calibration of laser-induced fluorescence excitation spectra of jet-cooled free radicals.

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